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## Identifying eccentric black hole mergers in dynamical formation environments

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#### Motivation: different formation channels





Active galactic nuclei



#### Motivation: the dynamical channel



How can we probe the dynamical origin of eccentric BBHs from *individual GW events?* 



## Setup



Fig. 2: setup of our system of interest.



Fig. 3: perturbed waveform due to Rømer delay.

$$
d\phi \approx 2\pi \frac{\Delta t(t)}{T_{12}(t)}
$$

#### Phase shift

 $\boxed{\Delta \phi (e) \approx \frac{288 \sqrt{2}}{85^2 g (1)^{13/2}} \frac{c^9}{G^{9/2}} \times \frac{m_3}{R^2} \frac{r_0^{13/2}}{m_1^2 m_2^2 m_{12}^{3/2}}}$  $\times e^{78/19}(1-e^2)^{1/2}g(e)^{13/2}$ 

#### Phase shift peaks when e~0.95!

#### How high can the phase shift get?

Take your favourite few-body simulation code

Simulate many binary-single interactions typical for clusters

Calculate the phase shift for those in triple configuration



#### How high can the phase shift get?



Fig. 4: maximum phase shift as a function of the peak GW frequency at formation of the binary. We use an equal mass triple of 20 solar masses. The different colours represent different initial semi-major axes: from 0.01 AU (blue) to 0.1 AU (red) and 1 AU (orange). The horizontal lines are the astrophysically expected upper limits.



How can we probe the dynamical origin of eccentric BBHs from *individual* GW events?





3G detectors such as LISA/ET/CE may be able to detect this environmental effect

# Questions?



Contact me at kai.hendriks@nbi.ku.dk!

#### Astrophysical scenarios: 3-body scatterings





Fig. 13: examples of 3-body scatterings that lead to phase shifts higher than the theoretical maximum.

#### Astrophysical scenario: 3-body scatterings



Fig. 8: 3-body scattering (5, 15 and 15 Msun) resulting in an inspiralling and merging binary on a bound orbit around a third object.

#### Astrophysical scenario: 3-body scatterings



Fig. 9: binary COM, reference trajectory and perturber trajectory of the scattering in question. We can extract the Rømer delay and phase shift from this.



Fig. 10: phase shift of this scattering, as a function of time Fig. 10: phase shift of this scattering, as a function of time  $12$ <br>(top) and binary eccentricity (bottom).

#### Phase shift: eccentric outer orbit



Fig. 7: phase shift as a function of time, for different realisations of  $f_{\rm m}$ , ranging from -¾  $\pi$ to ¾ π. We still use an eccentric 5 M $_{\odot}$  equal-mass binary, inspiralling on an eccentric orbit (e=0.9, a=30  $\rm R_{\odot}$ ) around a 100  $\rm M_{\odot}$  perturber. The binary assembles at a semi-major axis a $_{0}$ =1.3  $\rm R_{\odot}$  and e $_{0}$  = 0.9999.

## Stereoscopic images



#### Phase shift: analytical approximation

● Problem: no analytical description when outer orbit is eccentric



Fig. 9: trajectories and phase shift of an eccentric 5  $\rm M_{\odot}$  equal-mass binary, inspiralling on an eccentric orbit (e=0.9, a=30  $\rm R_{\odot}$ ) around a 100 M<sub> $_{\odot}$ </sub> perturber. The binaries assemble at a semi-major axis a<sub>0</sub>=1.3  $\rm R_{\odot}$  and e $_{\rm 0}$  = 0.9999, for 4 different outer phases at merger.

$$
\Delta\phi(e) \approx \frac{288\sqrt{2}}{85^2 g(1)^{13/2}} \frac{c^9}{G^{9/2}} \times \frac{m_3}{R^2} \frac{r_0^{13/2}}{m_1^2 m_2^2 m_{12}^{3/2}}
$$

$$
\times e^{78/19} (1 - e^2)^{1/2} g(e)^{13/2}
$$

#### Phase shift: circular outer orbit



Fig. 6: trajectory and phase shift of an eccentric 5 M<sub>☉</sub> equal-mass binary, inspiralling on a circular orbit (30 R<sub>☉</sub>) around a 100 M<sub>☉</sub> perturber. The binary assembles at a semi-major axis a $_{0}$ =1.3  ${\rm R}_{_{\bigodot}}$  and  ${\rm e}_{_{0}}$  :  $= 0.9999.$  16

#### Phase shift: effect of outer eccentricity



Fig. 8: trajectory and phase shift of an eccentric 5 M<sub>☉</sub> equal-mass binary, inspiralling on an eccentric orbit (e=0.9, a=30 R<sub>☉</sub>) around a 100 M<sub>☉</sub> perturber. The binary the straty the strate of the strate of the strate o assembles at a semi-major axis a $_{0}$ =1.3  $\rm R_{\odot}$  and e $_{0}$  = 0.9999, and merges at pericentre ( $\rm f_{m}^{\,}$ =

#### Phase shift: eccentric outer orbit



Fig. 10: trajectory and phase shift of an eccentric 5 M<sub>☉</sub> equal-mass binary, inspiralling on an eccentric orbit (e=0.9, a=30 R<sub>☉</sub>) around a 100 M<sub>☉</sub> perturber. The binary assembles at a semi-major axis a $_{0}$ =1.3 R $_{\odot}$  and e $_{0}$  = 0.9999, and merges at  $\rm{f_{m}^{\,}}$  :  $= -34$  π. 18

#### Phase shift: effect of the outer eccentricity



Fig. 9: trajectory and phase shift of an eccentric 5 M<sub>☉</sub> equal-mass binary, inspiralling on an eccentric orbit (e=0.9, a=30 R<sub>☉</sub>) around a 100 M<sub>☉</sub> perturber. The binary assembles at a semi-major axis a $_{0}$ =1.3  $\rm R_{\odot}$  and e $_{0}$  = 0.9999, and merges at apocentre ( $\rm f_{m}^{\phantom{\dag}}$  $=$ π). 19

#### Astrophysical scenario: 3-body scatterings



Fig. 11: 3-body scatterings with GW mergers in a stellar cluster..