Full and repeated TDEs

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TDES



(Original image from Rees88. Snapshots produced by Toscani using PHANTOM (Price+18)

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Stars tidally disrupted by BH tides (Rees88, Phinney89. Recent:: Ryu+20, Rossi+21, Bonnerot & Stone 21)

Multimessenger emitters

- Luminous electromagnetic flares (reviews: Saxton+20, vanVelzen+20, Alexander+20)
- High energy neutrinos

(Hayasaki21, Stein22, Reusch+22,)

(Kobayashi+04, **Toscani**+19,20,22,23, • GWs Pfister, **Toscani**+22)



TDEs



(Original image from Rees88. Snapshots produced by Toscani using PHANTOM (Price+18)

• Refine estimates of Pfister, Toscani+22 and Toscani+20 considering all the harmonics of the signal (formalism by Berry&Gair10)

 Distinction between full TDEs and repeated partial TDEs

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<u>DEs: main sequence stars</u>



 $\frac{d^4 \dot{N}_{\rm det}^{\rm fTDEs}}{dz dM_{\bullet} dM_{\ast} dR_{\rm p}} \approx \frac{d^2 \Gamma(M_{\bullet})}{dM_{\ast} dR_{\rm p}} \times \Phi(M_{\bullet}, z) \times \frac{4\pi \chi^2(z)}{H(z)} \times \Theta(z, M_{\bullet}, M_{\ast}, R_{\rm p})$

<u>DEs: main sequence stars</u>





(Hogg 1999)

<u>fTDEs: main sequence stars</u>





Chang et al. 2024

Broggi et al. 2024



<u>DEs: main sequence stars</u>



$\Theta(z, M_{\bullet}, M_{*}, R_{\rm p})$

 $\frac{d^4 \dot{N}_{\rm det}^{\rm fTDEs}}{dz dM_{\bullet} dM_{\ast} dR_{\rm p}} \approx \frac{d^2 \Gamma(M_{\bullet})}{dM_{\ast} dR_{\rm p}} \times \Phi(M_{\bullet}, z) \times \frac{4\pi \chi^2(z)}{H(z)} \times \Theta(z, M_{\bullet}, M_{\ast}, R_{\rm p})$

$$SNR^{2} = 4 \int_{0}^{\infty} df \frac{\tilde{h}(f)^{2}}{Sn_{n}(f)}$$
Harmonics

Instrument

<u>fTDEs: main sequence stars</u>



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$0.1 M_{\odot} \le M_* \le 10 M_{\odot}, 10^4 M_{\odot} \le M_{\bullet} \le 10^7 M_{\odot}, 10^{-4} \le z \le 3$

<u>fTDEs: white dwarfs</u>





Detected WD-TDEs in dwarf galaxies in 10 yrs

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$M_* = 0.5 M_{\odot}, 10^3 \le M_{\bullet} \le 10^5 M_{\odot}, 10^{-4} \le z \le 3$

<u>GW background from fTDEs</u>







10¹

rpTDEs: background

$$h_{\rm c,pop}^2 = \frac{G}{c^3 \pi^2} \times \frac{1}{f} \times \int_0^\infty dz \frac{d\dot{N}^{\rm tde}}{dz} \frac{1}{\chi^2(z)} \frac{1}{1+z} \left(\frac{dE}{df}\right)$$

Population of 1Msun stars disrupted by $10^5 - 3 \times 10^5 - 4 \times 10^6$ M_{\odot} (Broggi et al. 2024)

- two-body scattering
- number orbits per energy bin
- pericenter for total disruption

$$\frac{d\dot{N}}{dz} = \int d\epsilon \frac{d\dot{N}}{dzd\epsilon} = \int d\epsilon \frac{d\dot{N}}{d\epsilon} \frac{4\pi c\chi^2(z)}{H_0} \Phi(z)$$

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$$\frac{dE}{df} = \sum_{i=1}^{N_{\text{orbs}}(\Delta \epsilon_j)} \frac{dE}{df} \bigg|_{i}$$

<u>rpTDEs: background</u>



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rpTDEs: individual detections

 $P > \Delta t_{\rm obs}$



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rpTDEs: individual detections

 $P > \Delta t_{\rm obs}$





Conclusions

- fTDEs very unlikely sources for LISA, but promising for DO
- rpTDE background can be a problem for detectors more sensitive at lower frequency

• rpTDEs could be individually detected? LISA? DO?

QUESHONS?

<u>fTDEs: white dwarfs (II)</u>



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) ×
$$\Phi(M_{\bullet}, z)$$
 × $\frac{4\pi\chi^2(z)}{H(z)}$ × $\Theta(z, M_{\bullet}, M_*, R_p)$

$$N_{\rm gc} \approx \left(\frac{M_{\bullet}}{4 \times 10^5 {\rm M}_{\odot}}\right)$$

(Burkert & Tremaine10, Harris & Harris 11)

$$\begin{split} M_* &= 0.5 \,\mathrm{M}_\odot, M_\mathrm{h} = 10^3 \,\mathrm{M}_\odot, \\ 10^6 \,\mathrm{M}_\odot &\leq M_\bullet \leq 10^8 \,\mathrm{M}_\odot \end{split}$$

Repeated TDEs + full TDEs





Individual detections





$ilde{n}$: maximum number of subsequent passages

Empty loss cone

$\Delta t_{\rm obs} = \tilde{n}P + p$